

Branching Brownian motion with selection: a discrete model of front propagation

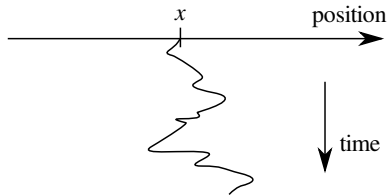
Pascal Maillard
(Université Paris-Sud)

SPA 2015, Oxford, July 13, 2015

Branching Brownian motion (BBM)

Definition

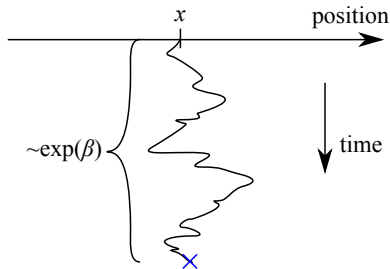
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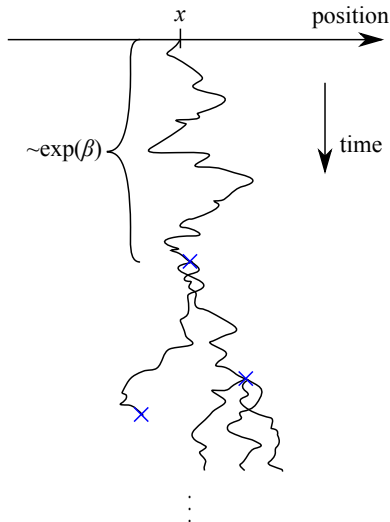
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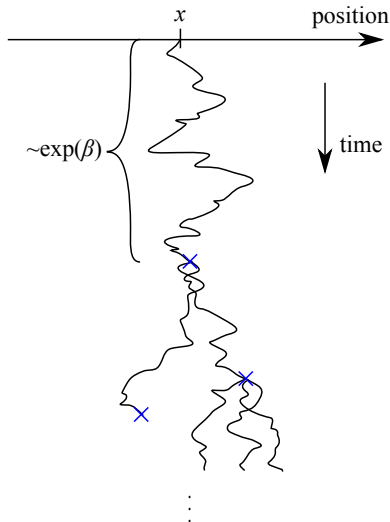


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→ A **Brownian motion** indexed by a **tree**.



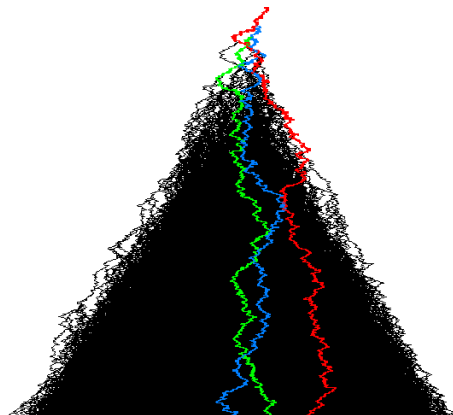
Branching Brownian motion (BBM) (2)

We always suppose
 $m := \mathbb{E}[L] - 1 > 0$.

Right-most particle

Let R_t be the position of the right-most particle. Then, as $t \rightarrow \infty$, almost surely on the event of survival,

$$\frac{R_t}{t} \rightarrow \sqrt{2\beta m}.$$



Picture by [Éric Brunet](#)

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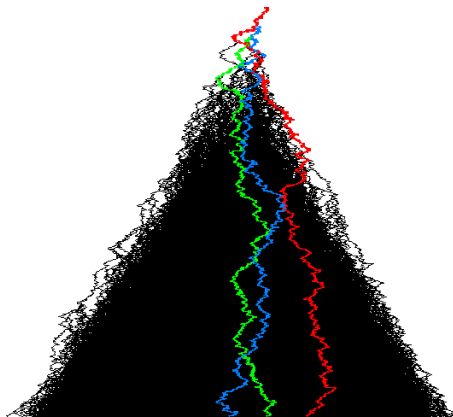
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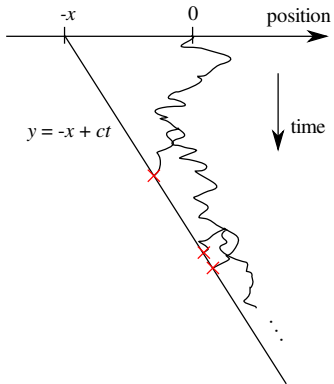
Convention

We will henceforth set
 $\beta = 1/(2m)$.

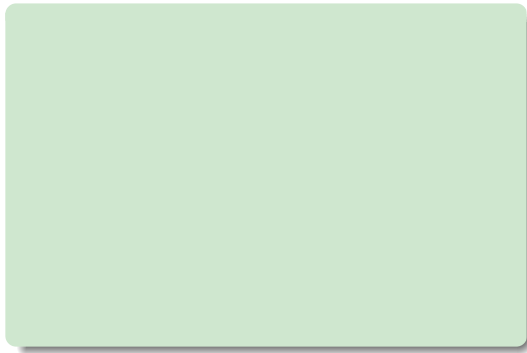


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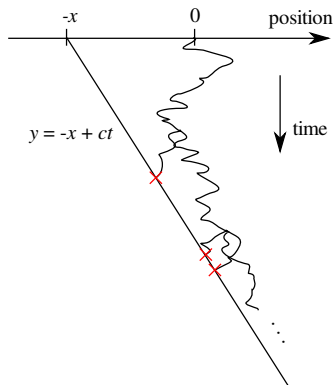
Selection



Two models of BBM with **selection**:



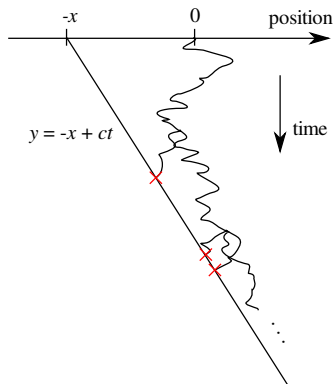
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- 1 **BBM with absorption:** Let $f(t)$ be a continuous function (the **barrier**). Kill an individual as soon as its position is less than $f(t)$.

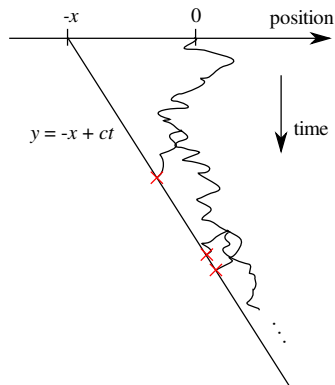
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Two models of BBM with **selection**:

- 1 **BBM with absorption:** Let $f(t)$ be a continuous function (the **barrier**). Kill an individual as soon as its position is less than $f(t)$.
- 2 **BBM with constant population size (N -BBM):** Fix $N \in \mathbb{N}$. As soon as the number of individuals exceeds N , only keep the N right-most individuals and kill the others.

Branching Brownian motion with absorption

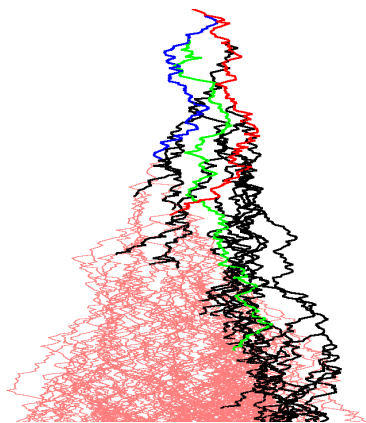


We take $f(t) = -x + ct$ (linear barrier).
Vast literature, known results (sample):

- almost sure extinction $\Leftrightarrow c \geq 1$
($c = 1$: **critical** case
 $c > 1$: **supercritical** case)
- growth rates for $c < 1$.
- asymptotics for extinction probability
for $c = 1 - \varepsilon$, ε small

Exact formulae for many quantities of interest.

BBM with constant population size

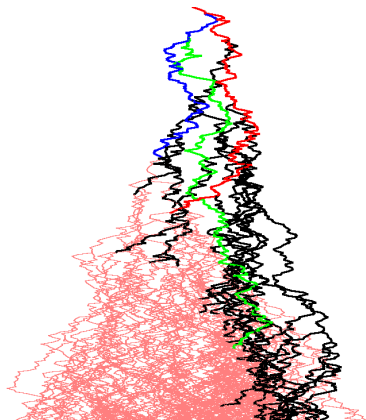


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Recall: Fix $N \in \mathbb{N}$. As soon as the number of individuals exceeds N , only keep the N right-most individuals and kill the others. **Much less tractable** than BBM with absorption:

- strong interaction between particles
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Nevertheless: A fairly detailed heuristic picture developed in the physics literature over the course of ten years:

Brunet and Derrida (1997-2004)
with Mueller and Munier (2006-2007)

- **Meta-stable** state: cloud of particles moving at **speed**

$v_N^{\text{det}} = \sqrt{1 - \pi^2 / \log^2 N}$, **empirical measure** seen from the left-most particle approximately proportional to $\sin(\pi x / \log N) e^{-x} \mathbf{1}_{(0, \log N)}(x)$,
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Real speed of the system is approximately

$$v_N = \sqrt{1 - \frac{\pi^2}{a_N^2}} = v_N^{\text{det}} + \frac{3\pi^2 \log \log N + o(1)}{\log^3 N},$$

and $O(1/\log^3 N)$ fluctuations.

Rigorous results (I)

- [Bérard and Guéré \(2012\)](#): Prove the $1/\log^2 N$ correction term to v_N for more general branching random walks (relying on results about BRW absorbed at a linear barrier by [Gantert, Hu and Shi \(2011\)](#)).
- [2×Berestycki and Schweinsberg \(2013\)](#): Study BBM with absorption at a linear barrier with slope $v_N \rightarrow$ toy model for N -BBM. They show convergence of the genealogy (as $N \rightarrow \infty$) to the Bolthausen–Sznitman coalescent, on time scale $(\log N)^3$.
- [Durrett and Remenik \(2010\)](#): Study empirical measure seen from left-most particle in a certain N -BRW. Show convergence of its evolution to a certain free-boundary convolution equation (*without rescaling in time*).
- [Mueller, Mytnik and Quastel \(2010\)](#): Prove $O(\log \log N / \log^3 N)$ correction term for noisy FKPP equation.

Rigorous results (2)

$$a_N = \log N + 3 \log \log N, \quad v_N = \sqrt{1 - \pi^2/a_N^2}$$

$X_i(t)$: position of i -th particle to the right at time t .

Theorem (M. '13+)

Suppose $\mathbb{E}[L^2] < \infty$ and at time 0, there are N particles distributed independently in $(0, a_N)$ according to density proportional to $\sin(\pi x/a_N)e^{-x}$. Then, for every $\alpha \in (0, 1)$,

$$(X_{\alpha N}(t \log^3 N) - v_N t \log^3 N)_{t \geq 0} \xrightarrow{\text{fidis}} (L_t)_{t \geq 0}.$$

Here, $(L_t)_{t \geq 0}$ is a Lévy process with $L_0 = x_\alpha$ (explicit), a certain (non-explicit) drift and explicit Lévy measure (the image of $\pi^2 x^{-2} \mathbf{1}_{x>0} dx$ by the map $x \mapsto \log(1+x)$).

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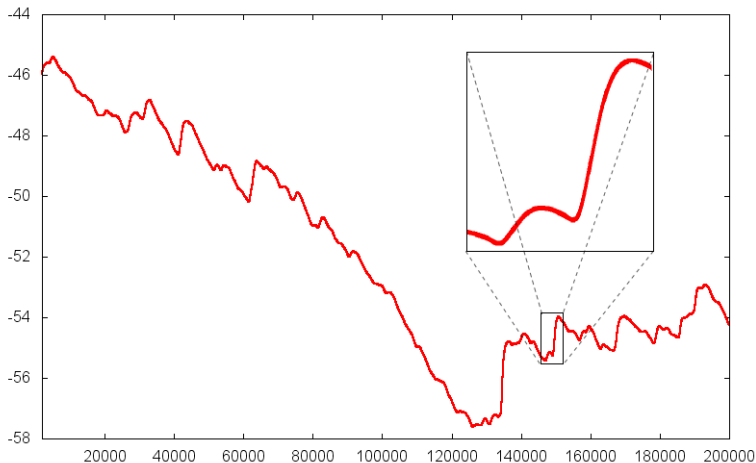
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Proof idea: Couple the N -BBM with BBM with a certain (random) absorbing barrier.

Simulation – Recentered position of barycenter



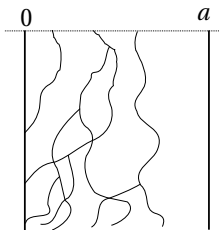
10^{10} particles

The B-BBM: the approximate model

$a = \log N + 3 \log \log N - A$: Position of a second barrier (idea from BBS '10).

Drift: $-\sqrt{1 - \pi^2/a^2}$.

Let first N , then A go to ∞ .



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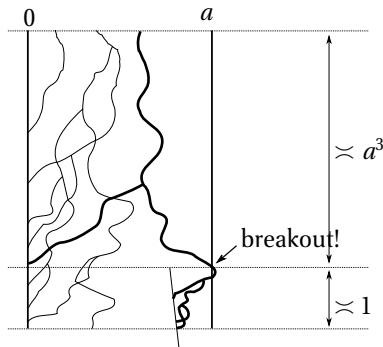
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When particle hits a , it will create $\asymp e^{-A} W N$ descendants, where (BBS '10)

$$\mathbb{P}(W > x) \sim x^{-1}, \quad x \rightarrow \infty.$$

Breakout when $W > \varepsilon e^A$, ε small.



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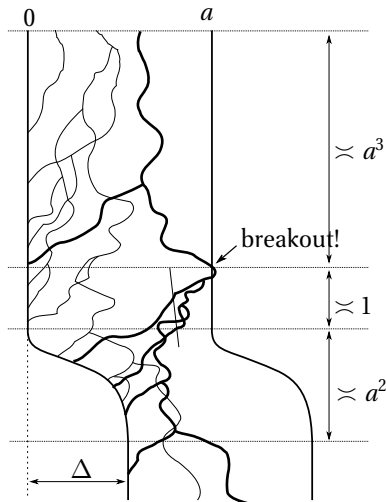
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After breakout, move barrier smoothly by random amount $\Delta \approx \log(1 + W)$.



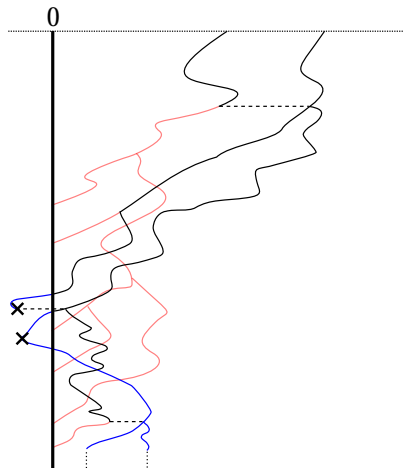
B-BBM \leftrightarrow N -BBM

First idea: couple both processes.

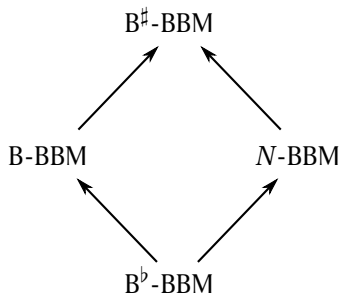
- **black** particles: present in B-BBM **and** N -BBM,
- **red** particles: present in B-BBM but **not** in N -BBM,
- **blue** particles: present in N -BBM but **not** in B-BBM.

Problem

Dependencies between particles too difficult to handle.



The solution



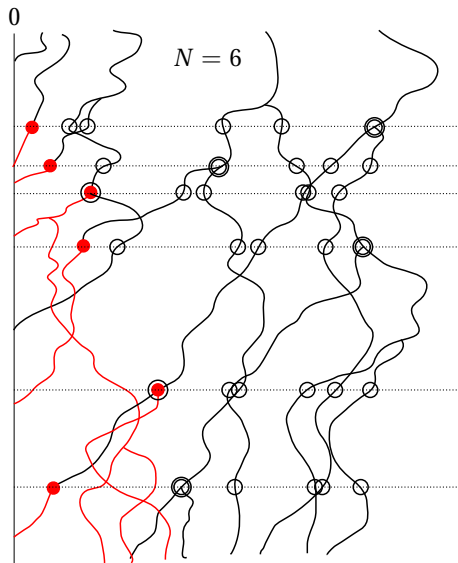
Introduce **two** auxiliary particle systems: The B^b -BBM and the $B^\#$ -BBM (stochastically) bound the N -BBM (and the B -BBM) from below and above (in the sense of stochastic order on the empirical measures).

Bounding the N -BBM from below: The B^b -BBM

Kill a particle

- whenever it hits 0 **or**
- whenever it has N particles to its right (red particles).

\implies more particles are being killed than in N -BBM.



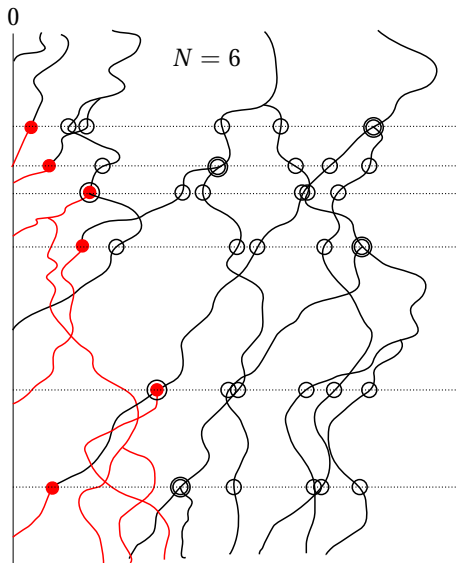
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At timescale $\log^3 N$, number of red particles stays negligible.

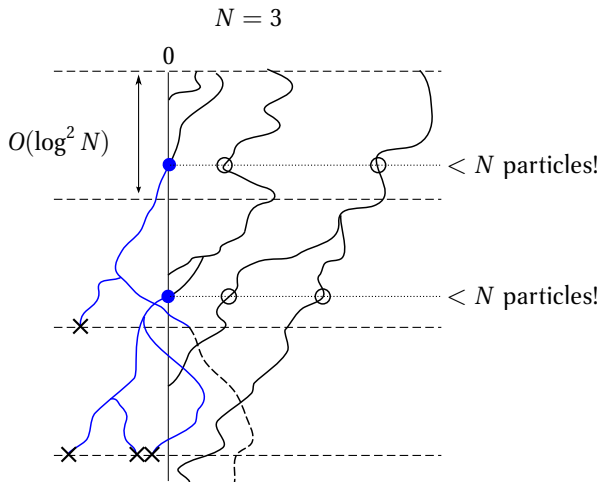


Bounding the N -BBM from above: The B^\sharp -BBM

Kill a particle whenever
it (at the same time)

- hits 0 **and**
- has N particles to its right.

A particle **survives**
temporarily
(blue particles) if it has
less than N particles to
its right the moment it
hits 0.



N -BBM open problems

- Long-time behavior:
 - exact speed asymptotics
 - empirical measure under equilibrium
 - relaxation time of empirical measure
- Genealogy
 - Show convergence to **Bolthausen–Sznitman coalescent** (at timescale $\log^3 N$). Proven for BBM with near-critical absorption (BBS '10) and for another N particle model called the *exponential model* (**Brunet–Derrida, Comets–Ramirez–Quastel, Cortines**)
- Durrett–Remenik free boundary equation
 - convergence to travelling wave

Related works/models

- N. Berestycki, Zhao '14: d -dimensional N -BBM (keep N particles with largest modulus). Show existence of a cloud of particles of width $\log N$ and length $(\log N)^{3/2}$ moving at linear speed in a uniformly chosen direction.
- Mallein '15: BBM (actually, branching random walk), fix $c > 0$. At time t , keep only $N_t = \exp(ct^{1/3})$ right-most particles (then $ct = (\log N)^3$). Position of right-most particle at time t :

$$t - \frac{3\pi^2}{2a^2}t^{1/3} + o(t^{1/3}).$$

Related works/models (2)

Mallein '15: BRW with slightly heavier tails ($\mathbb{E}[e^X] < \infty$ but possibly $\mathbb{E}[X^2 e^X] = \infty$) in the following regime:

- still linear speed of right-most particle
- path of right-most particle “almost” an excursion of an α -stable Lévy process, $\alpha \in (0, 2]$.

Considers N -BRW with these parameters. Shows that for some slowly varying function $L(x)$,

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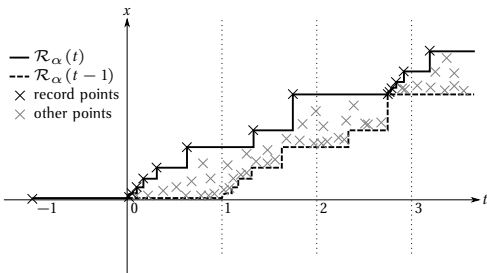
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Note: in all of these works, basic tool is coupling with BRW/BBM with absorption at a linear barrier.

Related works/models (3)

Bérard, M. '14: N -BRW with **regularly varying tails** (e.g. $\mathbb{P}(X > x) \sim x^{-\alpha}$, $\alpha > 0$). Binary branching. Phenomenology **much different than N -BBM**:

- Typically, most of the N particles are located near the minimum.
- From this position, **single** particles jump to higher positions and create new “colonies”.
- A colony reaches a population size of order N after time $\log_2 N$ (if it survived that long). At this time, it overtakes the whole population.
- For a colony to reach this size, it has to be created at a **new record position**.



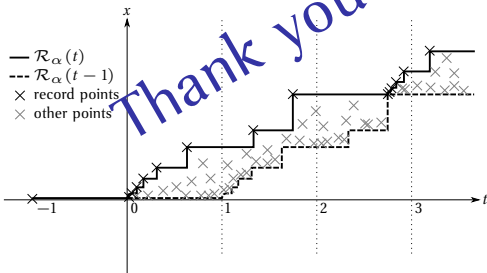
Limiting behavior described by a real-valued process $\mathcal{R}_\alpha(t)$ constructed out of a Poisson process. A consequence ($\alpha > 1$):

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Let $g : \mathbb{R} \rightarrow [0, 1]$ be measurable. Define

$$u(t, x) = \mathbb{E}_x \left[\prod_{u \in \mathcal{N}_t} g(X_u(t)) \right].$$

Then u satisfies the following partial differential equation:

Fisher–Kolmogorov–Petrovskii–Piskunov (FKPP) equation

$$\begin{cases} \partial_t u = \frac{1}{2} \partial_x^2 u + \beta(\mathbb{E}[u^L] - u) \\ u(0, x) = g(x) \quad (\text{initial condition}) \end{cases}$$

The prototype of a parabolic PDE admitting travelling wave solutions.

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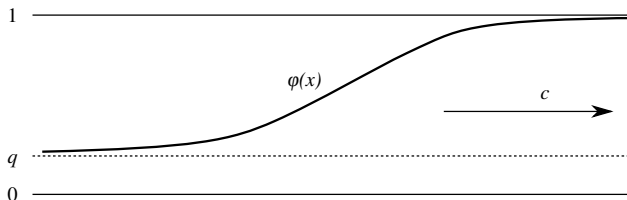
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Duality between BBM and FKPP.

Travelling waves



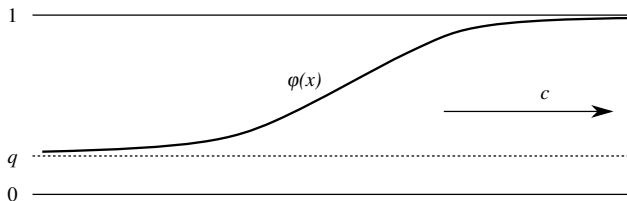
Definition

A *travelling wave* of speed c is a solution of the FKPP equation of the form

$$u(t, x) = \phi(x - ct),$$

where $\phi(x)$ is an increasing function with $\phi(\infty) = 1$ and $\phi(-\infty) = q$, where q solves $\mathbb{E}[q^L] = q$.

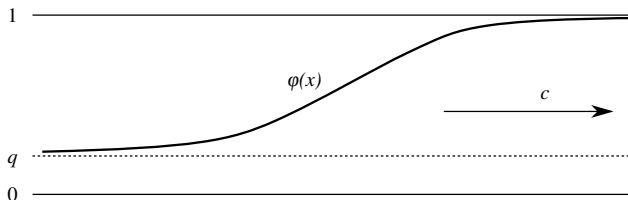
Travelling waves



Theorem (KPP '37)

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Travelling waves



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- Travelling waves *exist* for every speed $c \geq 1$ and are *unique* up to translation.
- Starting from Heaviside initial data $u(0, x) = 1_{\{x \geq 0\}}$, there exists a centering term $m(t)$, such that

$$u(t, x + m(t)) \xrightarrow{t \rightarrow \infty} \phi_1(x).$$

Noisy FKPP equation

$$\begin{cases} u(t, x) : \mathbb{R}_+ \times \mathbb{R} \rightarrow [0, 1] \\ \partial_t u = \partial_x^2 u + u(1-u) + \sqrt{\varepsilon u(1-u)} \dot{W} \\ u(0, x) = \mathbf{1}_{(x < 0)} \quad (\text{IC}) \end{cases}$$

Noisy FKPP equation

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- Admits travelling wave solutions with same phenomenology as N -BBM ($N \simeq \varepsilon^{-1}$) Mueller, Mytnik and Quastel '10